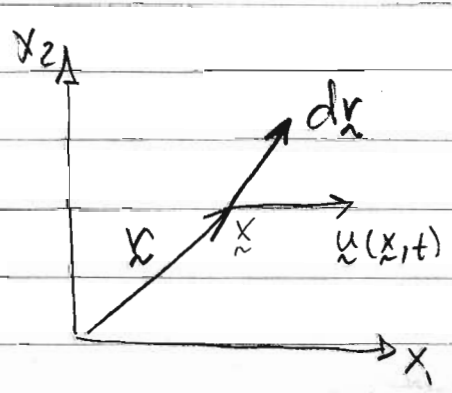


The relative motion near a point.

- Will give mathematical description of how small fluid elements deform in a given flow field.
- Physical understanding of mathematical expressions describing these deformations.
- Rates of deformation of fluid element directly linked to internal viscous forces in the fluid.



consider material line element

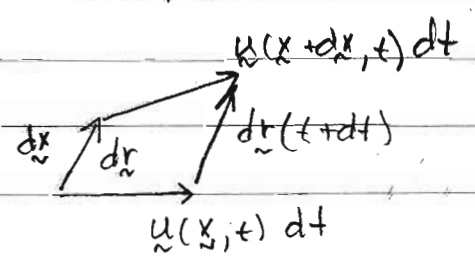
$$dr = r(x + dx, t) - r(x, t)$$

Translation with velocity

$$\frac{\partial r(x^0, \hat{t})}{\partial \hat{t}} = u(x, t); \quad x = r(x^0, \hat{t}), \quad t = \hat{t}$$

The instantaneous rate of change of length and orientation of a material line element is given by the relative motion.

$$\frac{\partial r(x^0 + dx^0, \hat{t})}{\partial \hat{t}} = u(x + dx, t); \quad dx = dr(x^0, dx^0, \hat{t}), \quad t = \hat{t}$$



$$\frac{\partial dr}{\partial \hat{t}} = \frac{D dr}{Dt} = u(x + dx, t) - u(x, t) = du$$

Taylor expansion around  $\underline{x} = 0$

$$\begin{aligned}
 du_i &= u_i(\underline{x} + d\underline{x}, t) - u_i(\underline{x}, t) = \\
 \text{Relative} &= u_i(\underline{x}, t) + \frac{\partial u_i(\underline{x}, t)}{\partial x_1} dx_1 + \frac{\partial u_i(\underline{x}, t)}{\partial x_2} dx_2 + \frac{\partial u_i(\underline{x}, t)}{\partial x_3} dx_3 + \\
 \text{velocity} &- u_i(\underline{x}, t) = \sum_{j=1}^3 \frac{\partial u_i(\underline{x}, t)}{\partial x_j} dx_j = \frac{\partial u_i(\underline{x}, t)}{\partial x_j} \underbrace{dx_j}
 \end{aligned}$$

$du_i$  are components of relative velocity at  $\underline{x}, t$

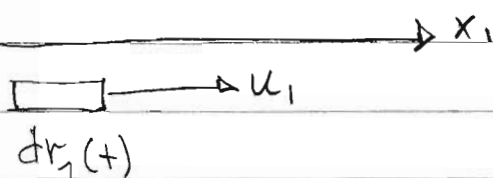
$dx_j$  are components of considered line element

$\frac{\partial u_i(\underline{x}, t)}{\partial x_j}$  the velocity gradient tensor  
(indep. of length & orient. of line element)

$$\frac{\partial u_i}{\partial x_j} = \begin{array}{|c|c|c|c|} \hline \frac{\partial u_1}{\partial x_1} & \frac{\partial u_1}{\partial x_2} & \frac{\partial u_1}{\partial x_3} & \\ \hline \frac{\partial u_2}{\partial x_1} & \frac{\partial u_2}{\partial x_2} & \frac{\partial u_2}{\partial x_3} & \\ \hline \frac{\partial u_3}{\partial x_1} & \frac{\partial u_3}{\partial x_2} & \frac{\partial u_3}{\partial x_3} & \\ \hline \end{array}$$

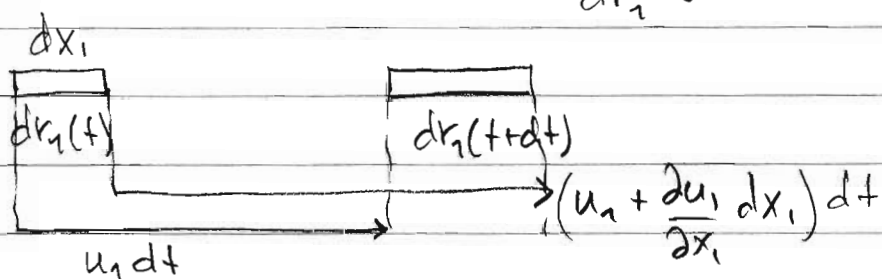
Evaluation of the effect on the relative motion from individual components of  $\frac{\partial u_i}{\partial x_j}$ .

\* Line element along  $x_1$ -axis



$$\frac{D}{Dt} dr_1 = \frac{\partial u_1}{\partial x_1} dr_1 + 0 + 0$$

Normal (linear) strain rate  $\left| \frac{D}{Dt} \frac{dr_1}{dr_1} = \frac{\partial u_1}{\partial x_1} \right.$

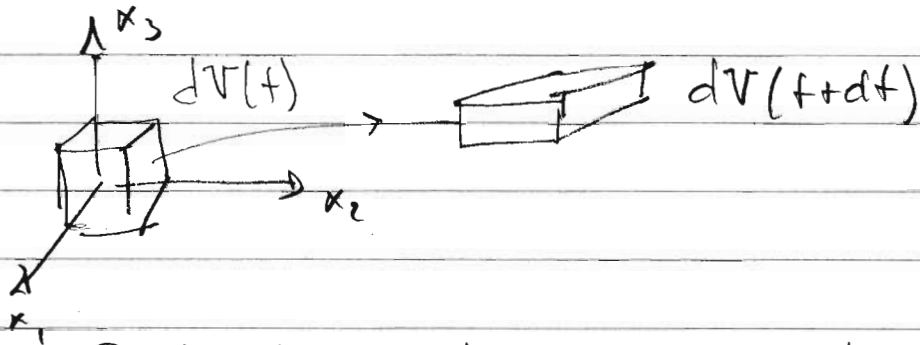


$$u_1 dt + dr_1(t+dt) = dr_1(t) + \left( u_1 + \frac{\partial u_1}{\partial x_1} dx_1 \right) dt$$

$$\frac{dr_1(t+dt) - dr_1(t)}{dt} = \frac{\partial u_1}{\partial x_1} dx_1$$

$$\frac{D}{Dt} \frac{dr_1}{dr_1}$$

\* Material volume  $dV = dr_1 dr_2 dr_3$



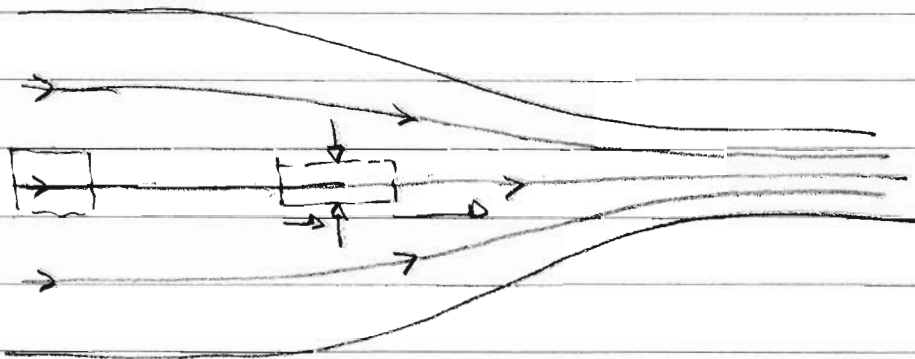
$$\text{Bulk strain rate } \frac{1}{dV} \frac{D dV}{Dt} = \frac{1}{dr_1 dr_2 dr_3} \frac{D(dr_1 dr_2 dr_3)}{Dt}$$

$$= \frac{1}{dr_1} \frac{D dr_1}{Dt} + \frac{1}{dr_2} \frac{D dr_2}{Dt} + \frac{1}{dr_3} \frac{D dr_3}{Dt} =$$

$$= \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} = \frac{\partial u_i}{\partial x_i} = \nabla \cdot \underline{u}$$

For incompressible fluid  $\frac{D dV}{Dt} = 0 \Leftrightarrow \nabla \cdot \underline{u} = 0$

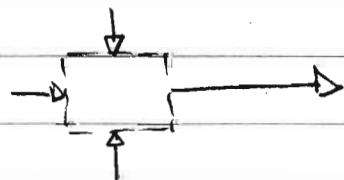
2D ex



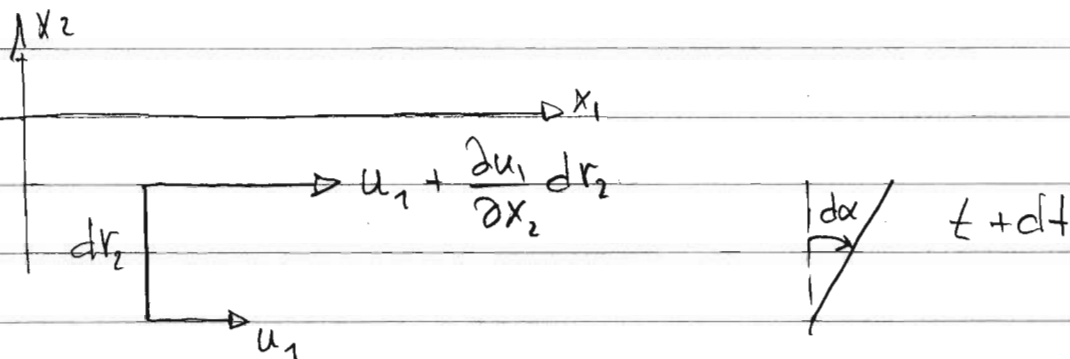
Fix control volume

$$\frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} = 0$$

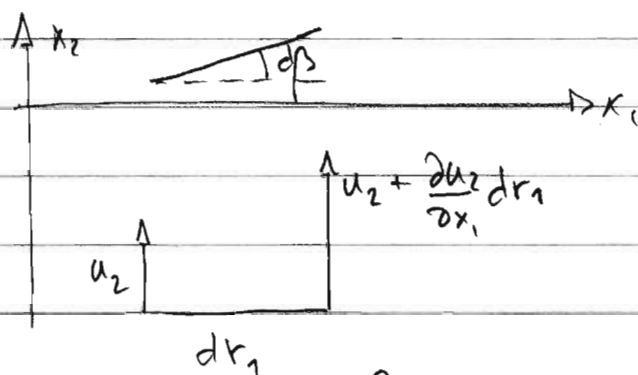
$$\underbrace{\quad}_{>0} \quad \underbrace{\quad}_{<0}$$



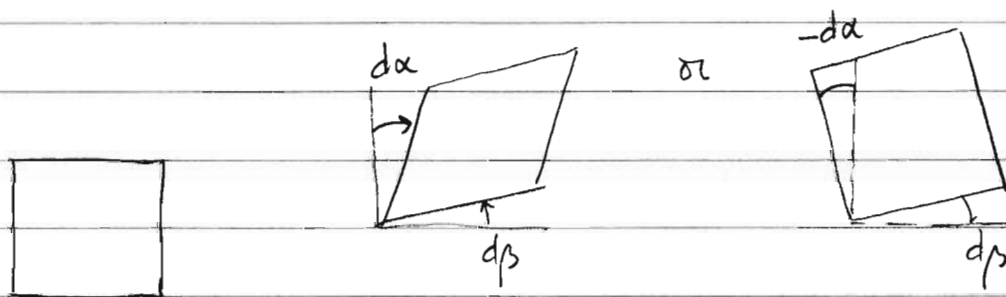
\* Line element  $\perp$  to the flow



$$\underbrace{\tan(\alpha)}_{\approx d\alpha} = \frac{dt \frac{\partial u_1}{\partial x_2} dr_2}{dr_2} \Rightarrow \frac{d\alpha}{dt} = \frac{\partial u_1}{\partial x_2}$$



$$\tan(\beta) = \frac{dt \frac{\partial u_2}{\partial x_1} dr_1}{dr_1} \Rightarrow \frac{d\beta}{dt} = \frac{\partial u_2}{\partial x_1}$$



Definitions: Shear strain rate

$$\frac{d\alpha}{dt} + \frac{d\beta}{dt} = \frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1}$$

gives deformation in shape

if  $d\beta = -d\alpha \Rightarrow$  no shape change, just rotation

## Vorticity

$$-\frac{dv_x}{dt} + \frac{dv_y}{dt} = \frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2} = \omega_3$$

$$\text{The average rotation rate} = \frac{1}{2} \omega_3 = \frac{1}{2} \left( \frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2} \right)$$

## Velocity gradient tensor

$$\frac{\partial u_i}{\partial x_j} = \underbrace{\frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)}_{\text{symmetric part}} + \underbrace{\frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} - \frac{\partial u_j}{\partial x_i} \right)}_{\text{antisymmetric part}}$$

symmetric part

antisymmetric part

$$e_{ij} = e_{ji}$$

$$\xi_{ij} = -\xi_{ji}$$

strain rate tensor

rotation tensor

$$e_{ij} = \begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{1}{2} \left( \frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) & \frac{1}{2} \left( \frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right) \\ \frac{1}{2} \left( \frac{\partial u_2}{\partial x_1} + \frac{\partial u_1}{\partial x_2} \right) & \frac{\partial u_2}{\partial x_2} & \frac{1}{2} \left( \frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right) \\ \frac{1}{2} \left( \frac{\partial u_3}{\partial x_1} + \frac{\partial u_1}{\partial x_3} \right) & \frac{1}{2} \left( \frac{\partial u_3}{\partial x_2} + \frac{\partial u_2}{\partial x_3} \right) & \frac{\partial u_3}{\partial x_3} \end{bmatrix}$$

$$e_{kk} = \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} = \nabla \cdot \underline{u}$$

$$\Xi_{ij} = \begin{bmatrix} 0 & -\frac{1}{2}\omega_3 & \frac{1}{2}\omega_2 \\ \frac{1}{2}\omega_3 & 0 & -\frac{1}{2}\omega_1 \\ -\frac{1}{2}\omega_2 & \frac{1}{2}\omega_1 & 0 \end{bmatrix}$$

$$\omega_3 = \frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2}$$

$$\omega_1 = \frac{\partial u_3}{\partial x_2} - \frac{\partial u_2}{\partial x_3}$$

$$\omega_2 = \frac{\partial u_1}{\partial x_3} - \frac{\partial u_3}{\partial x_1}$$

Shorter notation

$$\Xi_{ij} = -\frac{1}{2} \underbrace{\epsilon_{ijk}}_{\text{permutation tensor}} \omega_k$$

permutation tensor

$\epsilon_{ijk} = 1$  for cyclic permutations of 123

$\epsilon_{ijk} = -1$  for non-cyclic ——— " ———

$\epsilon_{ijk} = 0$  otherwise

Cross product  $(\underline{a} \times \underline{b})_i = \epsilon_{ijk} a_j b_k$

Rotation of 3D line element

$$\begin{aligned} \left( \frac{D}{Dt} \frac{d\underline{r}}{dt} \right)_{\text{rot.}} &= \Xi_{ij} dx_j = -\frac{1}{2} \epsilon_{ijk} \omega_k dx_j = -\frac{1}{2} \underline{dx} \times \underline{\omega} \\ &= \frac{1}{2} \underline{\omega} \times \underline{dx} \end{aligned}$$

i.e. solid body rotation at angular velocity  $\frac{1}{2}\underline{\omega}$   
(c.f.  $\dot{\underline{r}} = \underline{\Omega} \times \underline{r}$ )

One also finds that

$$\vec{\omega} = \nabla \times \vec{u} = \begin{vmatrix} \vec{e}_1 & \vec{e}_2 & \vec{e}_3 \\ \frac{\partial}{\partial x_1} & \frac{\partial}{\partial x_2} & \frac{\partial}{\partial x_3} \\ u_1 & u_2 & u_3 \end{vmatrix}$$

Tensor notation  $\omega_i = \epsilon_{ijk} \frac{\partial u_k}{\partial x_j} =$

$$= \underbrace{\epsilon_{ijk} e_{kj}}_{\substack{=0 \\ \text{as } \epsilon}} + \underbrace{\epsilon_{ijk} \zeta_{kj}}$$

$$\left. \begin{array}{l} \begin{bmatrix} a & c \\ c & b \end{bmatrix} \begin{bmatrix} 0 & d \\ -d & 0 \end{bmatrix} = 0 + cd - cd + 0 = 0 \\ s_{ij} \quad a_{ij} \end{array} \right\}$$

$$\begin{aligned} \left[ -\frac{1}{2} \epsilon_{ijk} \omega_k &= -\frac{1}{2} \epsilon_{ijk} \epsilon_{klem} \zeta_{me} = -\frac{1}{2} \epsilon_{kij} \epsilon_{klem} \zeta_{me} \right. \\ &= -\frac{1}{2} \zeta_{me} (\delta_{ie} \delta_{jm} - \delta_{im} \delta_{je}) = \\ &= -\frac{1}{2} \zeta_{ji} + \frac{1}{2} \zeta_{ij} = \frac{1}{2} \zeta_{ij} + \frac{1}{2} \zeta_{ij} = \zeta_{ij} \end{aligned}$$



## Partition of $e_{ij}$

- isotropic part

$$\bar{e}_{ij} = \begin{bmatrix} \frac{1}{3} e_{kk} & 0 & 0 \\ 0 & \frac{1}{3} e_{kk} & 0 \\ 0 & 0 & \frac{1}{3} e_{kk} \end{bmatrix} = \frac{1}{3} e_{kk} \delta_{ij} = \frac{1}{3} \nabla \cdot \mathbf{u} \delta_{ij}$$

gives deformation due to volume expansion

- deviatoric part

$$\bar{e}_{ij} = e_{ij} - \bar{e}_{ij} = \frac{1}{2} \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} - \frac{2}{3} \frac{\partial u_k}{\partial x_k} \delta_{ij} \right)$$

gives deformation rate without change of volume.

$$\text{traceless } \bar{e}_{kk} = 0$$

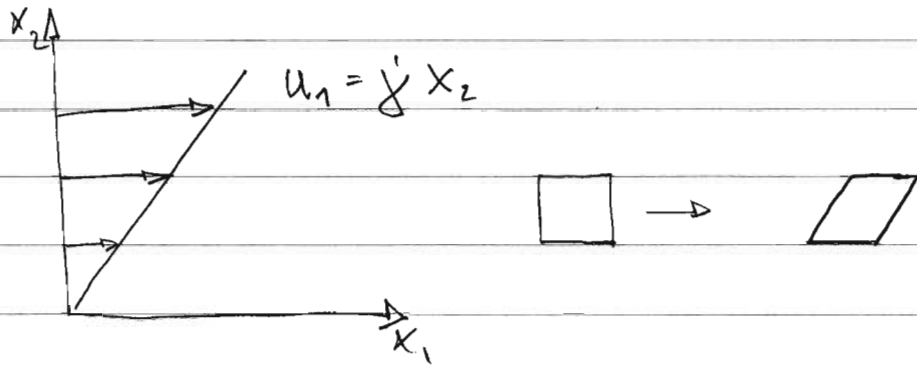
## Summary of relative motion

$$\frac{D}{Dt} dr_i = du_i = \underbrace{\frac{\partial u_i}{\partial x_k} dx_k}_{(1)} + \underbrace{\bar{e}_{ik} dx_k}_{(2)} + \underbrace{\bar{\omega}_{ik} dx_k}_{(3)}$$

The terms of  $du_i$  are results of

- (1) pure deformation, without change of volume
- (2) isotropic change of volume
- (3) local solid body rotation

Ex: Shear flow

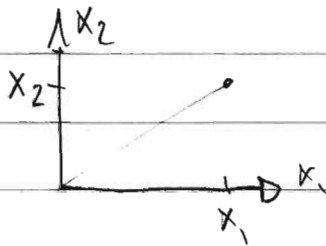


$$\frac{\partial u_i}{\partial x_j} = \begin{bmatrix} 0 & \dot{\gamma} & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} = \underbrace{\begin{bmatrix} 0 & \frac{1}{2}\dot{\gamma} & 0 \\ \frac{1}{2}\dot{\gamma} & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}}_{\bar{e}_{ij}} + \underbrace{\begin{bmatrix} 0 & \frac{1}{2}\dot{\gamma} & 0 \\ -\frac{1}{2}\dot{\gamma} & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}}_{\bar{\zeta}_{ij}}$$

$$e_{nn} = \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} = 0 + 0 + 0 \quad \text{no vol. def.}$$

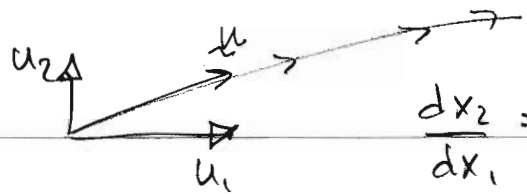
Isotropic strain rate tensor  $\bar{e}_{ij} = 0$

Relative velocity field  $du_i = \frac{\partial u_i}{\partial x_k} dx_k =$



$$= \frac{1}{2} \begin{bmatrix} 0 & \dot{\gamma} & 0 \\ \dot{\gamma} & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ 0 \end{bmatrix} + \frac{1}{2} \begin{bmatrix} 0 & \dot{\gamma} & 0 \\ -\dot{\gamma} & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ 0 \end{bmatrix} = \underbrace{\dot{\gamma} \begin{bmatrix} x_2 \\ x_1 \\ 0 \end{bmatrix}}_{(du_i)_{def.}} + \underbrace{\dot{\gamma} \begin{bmatrix} x_2 \\ -x_1 \\ 0 \end{bmatrix}}_{(du_i)_{ro}}$$

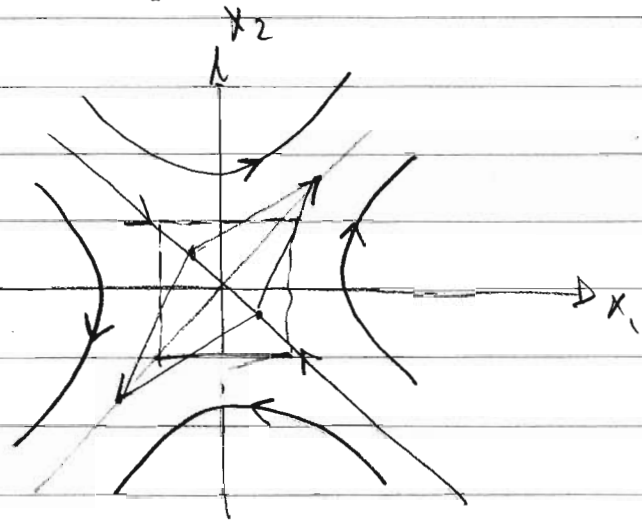
Streamlines



$$\frac{dx_2}{dx_1} = \frac{u_2}{u_1}$$

Deformation  $\frac{dx_2}{dx_1} = \frac{\frac{1}{2} \gamma x_1}{\frac{1}{2} \gamma x_2} = \frac{x_1}{x_2} \Rightarrow x_2 dx_2 = x_1 dx_1$

$$\Rightarrow \frac{x_2^2}{2} = \frac{x_1^2}{2} + c$$



Rotation  $\frac{dx_2}{dx_1} = \frac{-\frac{1}{2} \gamma x_1}{\frac{1}{2} \gamma x_2} = -\frac{x_1}{x_2} \Rightarrow x_2 dx_2 + x_1 dx_1 = 0$

$$\Rightarrow \frac{x_2^2}{2} + \frac{x_1^2}{2} = c$$

