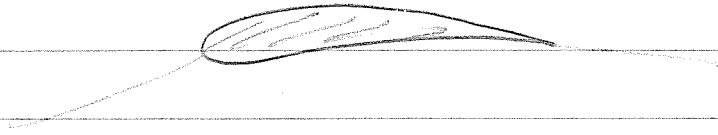
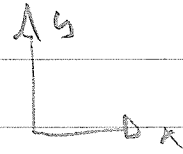


2 D irrotational flow.

$$\omega_z = 0$$



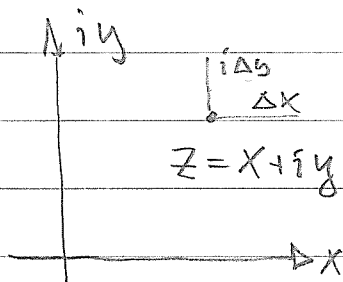
Velocity potential $u = \nabla \phi$

Incompressible \Rightarrow $\nabla^2 \phi = 0$

Stream function $u = \frac{\partial \psi}{\partial y}$; $v = -\frac{\partial \psi}{\partial x}$

$\omega_z = 0 \Rightarrow$ $\nabla^2 \psi = 0$

Method of solution with analytic functions.



Complex function $F(z)$
is analytic if $\frac{dF(z)}{dz}$

exists and is independent
of direction.

$$F'(z) = \lim_{\Delta z \rightarrow 0} \left(\frac{F(z + \Delta z) - F(z)}{\Delta z} \right)$$

$$\text{Let } F(z) = \phi(x, y) + i\psi(x, y)$$

$$F' = \lim_{\Delta x \rightarrow 0} \frac{1}{\Delta x} \left\{ \phi(x+\Delta x, y) + i\psi(x+\Delta x, y) - \phi(x, y) - i\psi(x, y) \right\} =$$

$$= \frac{\partial \phi}{\partial x} + i \frac{\partial \psi}{\partial x}$$

$$F' = \lim_{i\Delta y \rightarrow 0} \frac{1}{i\Delta y} \left\{ \phi(x, y+\Delta y) + i\psi(x, y+\Delta y) - \phi(x, y) - i\psi(x, y) \right\} =$$

$$= \frac{1}{i} \frac{\partial \phi}{\partial y} + \frac{\partial \psi}{\partial y} = \frac{\partial \psi}{\partial y} - i \frac{\partial \phi}{\partial y}$$

$$\left. \begin{array}{l} \text{Real part: } \frac{\partial \phi}{\partial x} = \frac{\partial \psi}{\partial y} \\ \text{Imaginary part: } \frac{\partial \psi}{\partial x} = -\frac{\partial \phi}{\partial y} \end{array} \right\} \text{Cauchy-Riemann equations}$$

One finds:

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = \frac{\partial^2 \psi}{\partial x \partial y} - \frac{\partial^2 \psi}{\partial y \partial x} = 0$$

$$\nabla^2 \psi = \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} = -\frac{\partial^2 \phi}{\partial x \partial y} + \frac{\partial^2 \phi}{\partial y \partial x} = 0$$

$$\nabla \phi \cdot \nabla \psi = \frac{\partial \phi}{\partial x} \frac{\partial \psi}{\partial x} + \frac{\partial \phi}{\partial y} \frac{\partial \psi}{\partial y} = -\frac{\partial \phi}{\partial x} \frac{\partial \phi}{\partial y} + \frac{\partial \phi}{\partial y} \frac{\partial \phi}{\partial x} = 0$$

iso level surfaces of ϕ & ψ are orthogonal.

Any analytic function $F(z) = \phi + i\psi$ is a candidate to describe an irrotational, incompressible 2D flow field with:

$$u = \frac{\partial \phi}{\partial x} = \frac{\partial \psi}{\partial y}, \quad v = \frac{\partial \phi}{\partial y} = -\frac{\partial \psi}{\partial x}$$

We define the complex velocity $W(z)$

$$\begin{aligned} W(z) &= \frac{dF}{dz} = \frac{\partial \phi}{\partial x} + i \frac{\partial \psi}{\partial x} = u - i v \\ &= \frac{\partial \psi}{\partial y} - i \frac{\partial \phi}{\partial x} = u - i v \end{aligned}$$

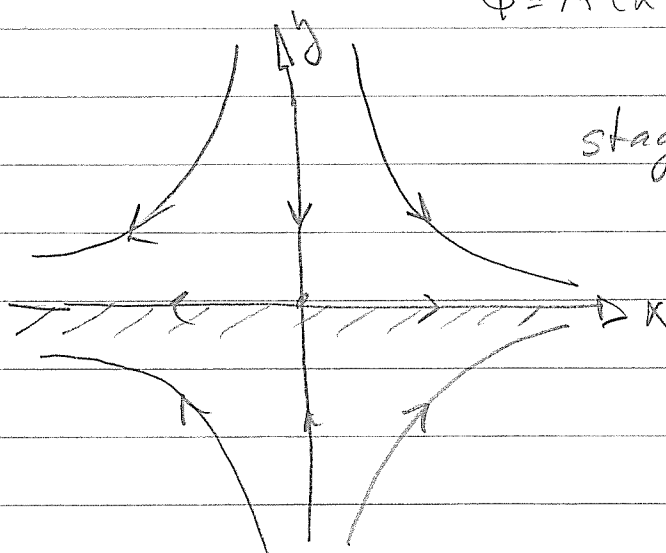
W is complex conjugate of velocity vector.

Ex. $F = A z^2 \quad W = 2Az = 2A(x+iy)$

$$\begin{cases} u = 2Ax \\ v = -2Ay \end{cases}$$

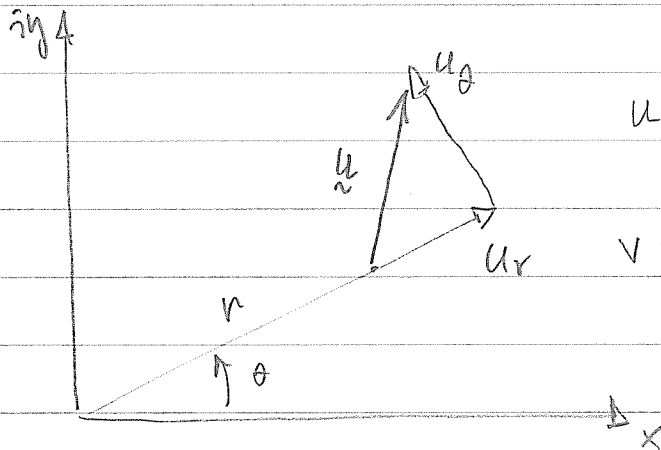
$$Az^2 = A(x+iy)^2 = A(x^2 - y^2) + 2Aixy$$

$$\phi = A(x^2 - y^2), \quad \psi = 2Axy$$



stagnation point flow

Polar coordinates $z = r e^{i\theta} = \underbrace{r \cos\theta}_x + i \underbrace{r \sin\theta}_y$



$$u = u_r \cos\theta - u_\theta \sin\theta$$

$$v = u_r \sin\theta + u_\theta \cos\theta$$

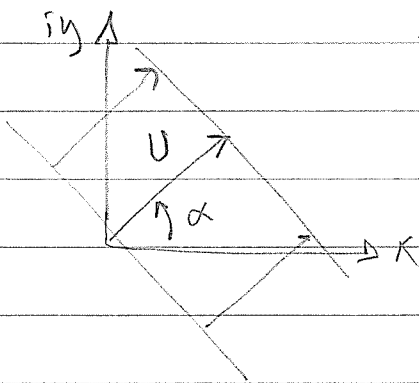
$$W = u - iv = \cos\theta (u_r - iu_\theta) + \sin\theta \underbrace{(-iu_r - u_\theta)}_{-i \sin\theta (u_r - iu_\theta)}$$

$$= (u_r - iu_\theta) (\cos\theta - i \sin\theta) =$$

$$= \underbrace{(u_r - iu_\theta)}_{-i\theta} e^{-i\theta}$$

Ex. $F = U e^{-i\alpha} z$

$$w = F' = U e^{-i\alpha} = \underbrace{U \cos\alpha}_u - i \underbrace{U \sin\alpha}_v$$



$$W = U e^{-i\alpha + i\theta - i\theta} = U e^{i(\theta - \alpha)} e^{-i\theta}$$

$$= \left(\underbrace{U \cos(\theta - \alpha)}_{u_r} + i \underbrace{U \sin(\theta - \alpha)}_{-u_\theta} \right) e^{-i\theta}$$

Ex. $F = A z^n$ (A real, $n \geq \frac{1}{2}$)

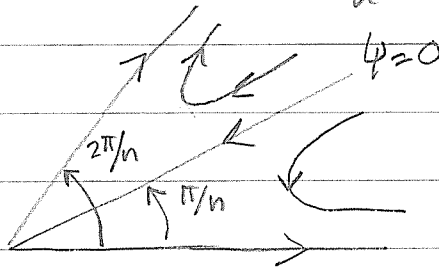
$$w = F' = n A z^{n-1} = n A r^{n-1} e^{i n \theta} e^{-i \theta}$$

$$= n A r^{n-1} (\cos(n\theta) + i \sin(n\theta)) e^{-i \theta}$$

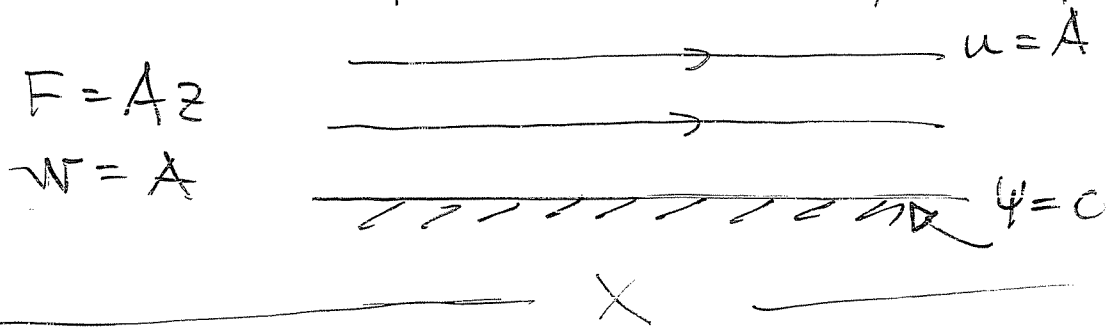
$$\Rightarrow \begin{cases} u_r = n A r^{n-1} \cos(n\theta) \\ u_\theta = -n A r^{n-1} \sin(n\theta) \end{cases}$$

$$F = A r^n e^{i n \theta} = \underbrace{A r^n \cos(n\theta)}_{\phi(r, \theta)} + i \underbrace{A r^n \sin(n\theta)}_{\psi(r, \theta)}$$

$$\psi = A r^n \sin(n\theta) = 0 \Rightarrow \theta = \frac{\pi k}{n} \quad k = 0, 1, 2, \dots$$



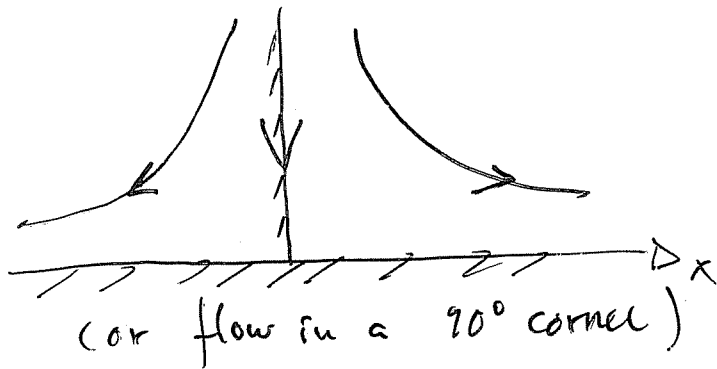
$n=1$: uniform flow over flat plate



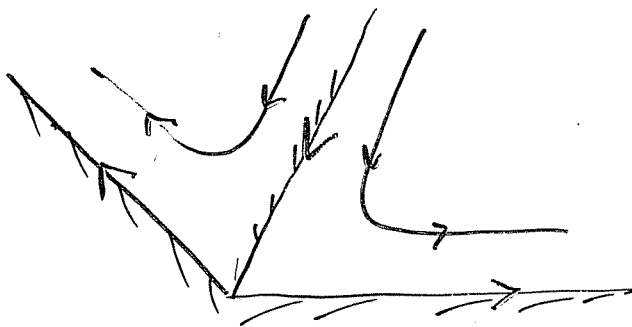
$n=2$: stagnation point flow

$F = Az^2$ $W = 2Az = 2A(x + iy)$

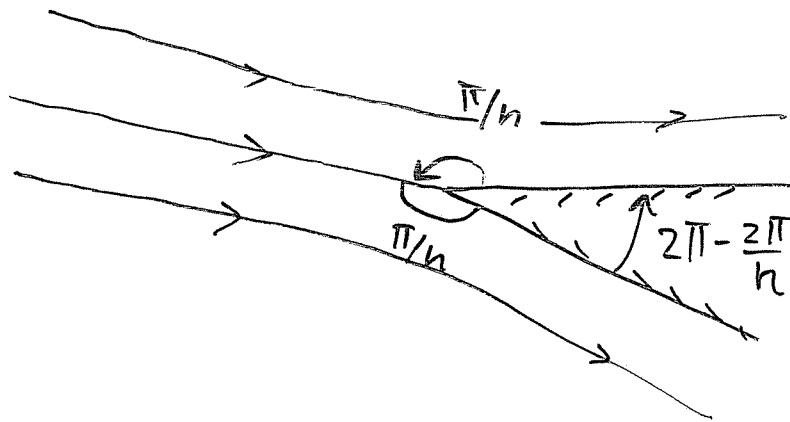
$\begin{cases} u = 2Ax \\ v = -2Ay \end{cases}$



$n > 2$:

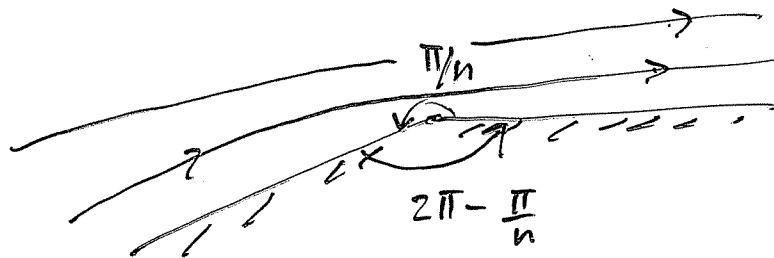


$1 \leq n \leq 2$: flow towards a wedge
(or in a corner)



at surface $\vartheta = 0$: $u_r = n A r^{n-1}$

$n < 1$: flow around corner



at surface $\vartheta = 0$: $u_r = \frac{n A}{r^{1-n}}$

$\vartheta = \frac{\pi}{n}$: $u_r = -\frac{n A}{r^{1-n}}$

infinite
velocity
at corner!

$n = \frac{1}{2}$: flow around edge



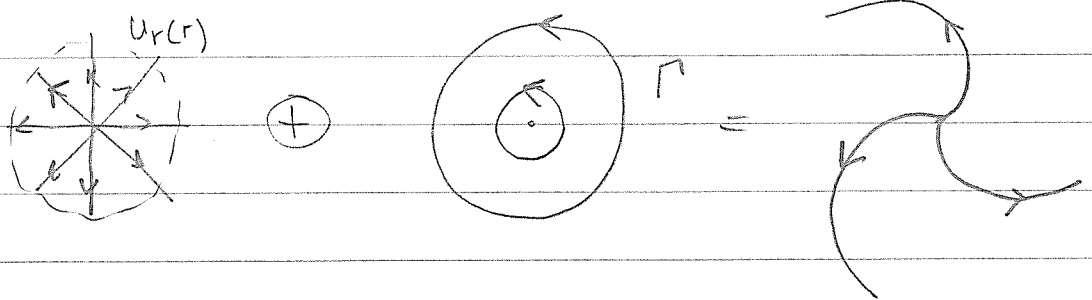
Ex. Line source & line vortex

$$F(z) = \frac{(m - i\Gamma)}{2\pi} \ln z \quad z = r e^{i\theta}$$

$$W = \frac{m - i\Gamma}{2\pi} \frac{1}{z} = \frac{(m - i\Gamma)}{2\pi r} e^{-i\theta} \triangleq (u_r - i u_\theta) e^{-i\theta}$$

$$u_r = \frac{m}{2\pi r} \quad m = 2\pi r u_r(r) \text{ is volume flux/width}$$

$$u_\theta = \frac{\Gamma}{2\pi r} \quad \Gamma \text{ is circulation}$$



$$F = \frac{(m - i\Gamma)}{2\pi} (\ln r + i\theta) \triangleq \phi + i\psi$$

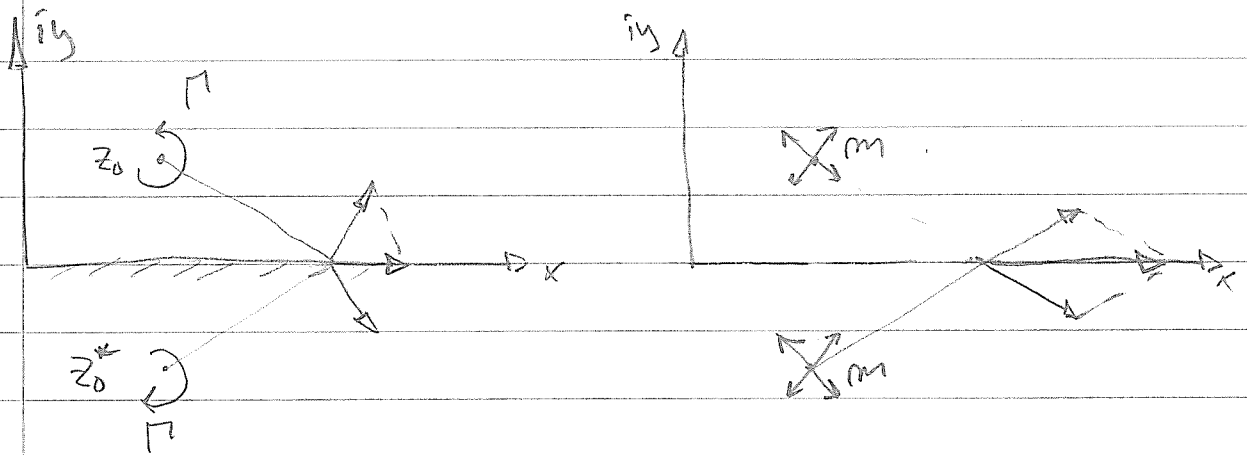
$$\phi = \frac{m}{2\pi} \ln r + \frac{\Gamma}{2\pi} \theta$$

$$\psi = \frac{m}{2\pi} \theta - \frac{\Gamma}{2\pi} \ln r$$

Superpositions of analytic functions give new analytic solution.

Laplace equation is linear!

Ex. Mirror images in a plane

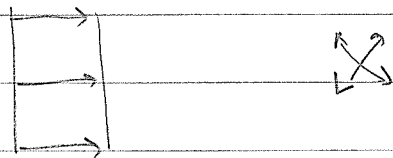


mirror images separated

$$F = \frac{(m - i\Gamma)}{2\pi} \ln(z - z_0) + \frac{(m + i\Gamma)}{2\pi} \ln(z - z_0^*)$$

Ex. Line source \oplus uniform stream

$$F = Uz + \frac{m}{2\pi} \ln z = Ure^{i\theta} + \frac{m}{2\pi} (\ln r + i\theta)$$

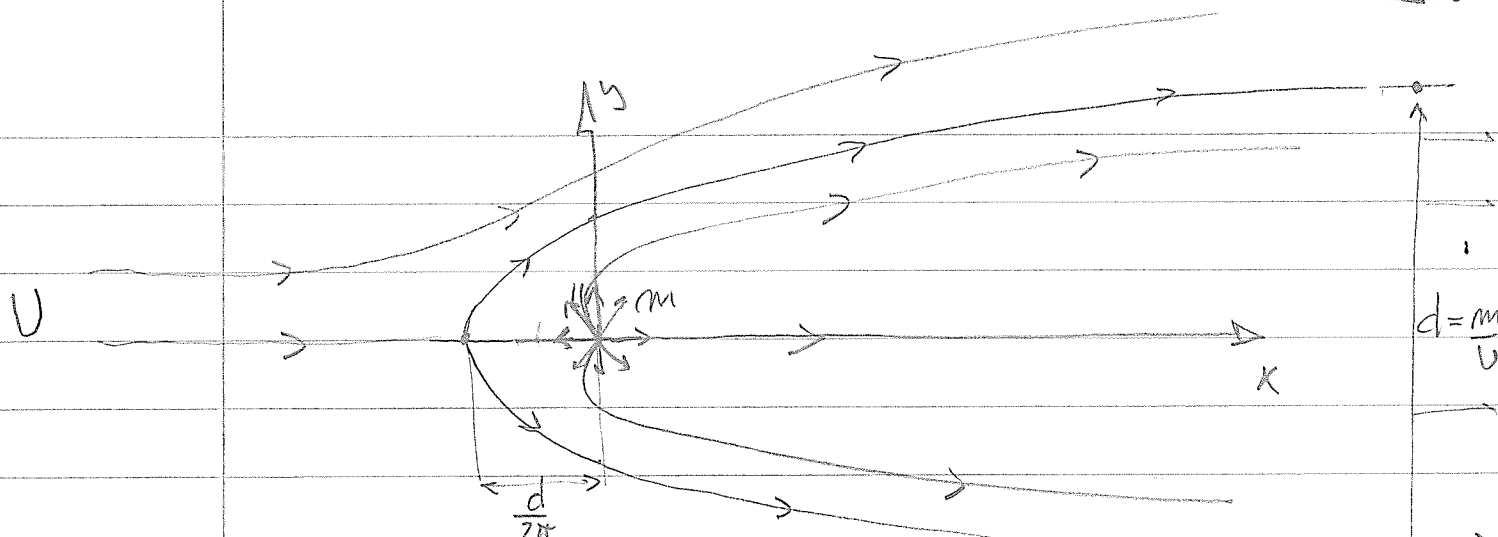


Stream function $\psi = Ur \sin\theta + \frac{m}{2\pi} \theta$

$$w = F' = U + \frac{m}{2\pi z}$$

Stagnation point if $w = 0$

$$U + \frac{m}{2\pi z_s} = 0 \Rightarrow z_s = -\frac{m}{2\pi U}$$



Streamline through stagnation point?

$$\psi = \text{constant} = ? \quad z_s = -\frac{m}{2\pi U} \Rightarrow r_s = \frac{m}{2\pi U}, \vartheta_s = \pi$$

$$\psi_s = 0 + \frac{m}{2} = U r_s \sin \vartheta + \frac{m}{2\pi} \vartheta = y$$

Parametric representation

$$y = \frac{1}{2} \frac{m}{U} - \frac{m}{U} \frac{\vartheta}{2\pi}$$

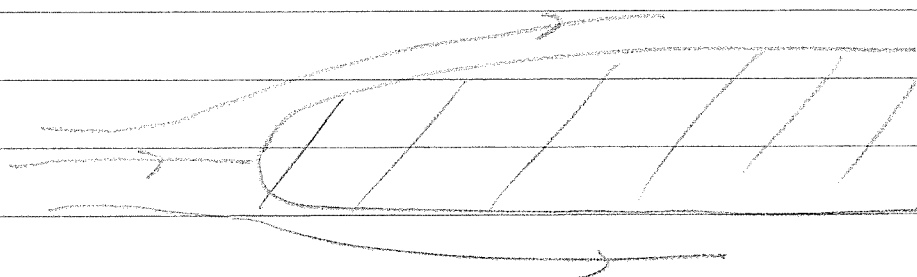
$$= d \left(\frac{1}{2} - \frac{\vartheta}{2\pi} \right)$$

$$\left. \begin{array}{l} \vartheta \rightarrow 0 \quad y \rightarrow \frac{1}{2} \frac{m}{U} \\ \vartheta \rightarrow 2\pi \quad y \rightarrow -\frac{1}{2} \frac{m}{U} \end{array} \right\} d = \frac{m}{U}$$

$$\vartheta = \pi \quad y = 0, \quad r_s = \frac{d}{2\pi}$$

$$\vartheta = \pm \pi/2, \quad y = \pm d/4$$

Flow past half-infinite body.



Ex. Dipole (doublet)

$$F = \frac{\mu}{2\pi z} \quad \mu \text{ is dipole strength}$$

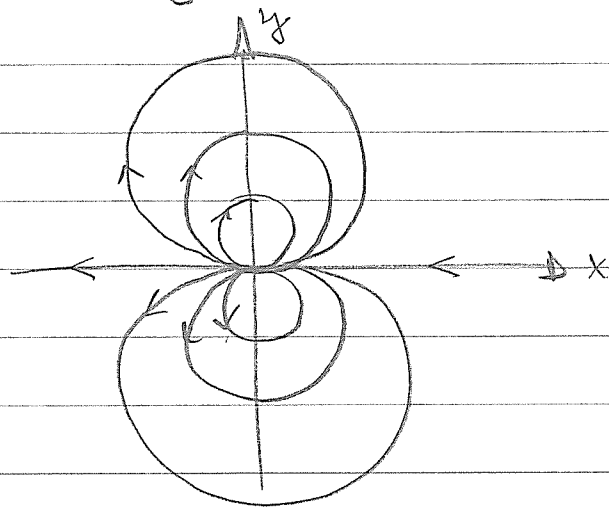
$$= \frac{\mu z^*}{2\pi z z^*} = \frac{\mu}{2\pi r^2} (x - iy)$$

$$= \underbrace{\frac{\mu x}{2\pi (x^2 + y^2)}}_{\Phi} - i \underbrace{\frac{\mu y}{2\pi (x^2 + y^2)}}_{\Psi}$$

$$\Psi = -\frac{\mu y}{2\pi (x^2 + y^2)} = C \Rightarrow \text{stream lines}$$

$$\frac{y}{x^2 + y^2} = \frac{1}{2R} \Rightarrow x^2 + y^2 - 2Ry = 0$$

$$x^2 + (y - R)^2 - R^2 = 0$$



One finds:

$$u_r = -\frac{\mu}{2\pi r^2} \cos \theta, \quad u_\theta = -\frac{\mu}{2\pi r^2} \sin \theta$$