

Lecture I.

{ Show Powerpoint presentation as introduction }
 ~ 30 pages probably ~ 45 min

Vector form & Cartesian tensor form of equations.

{ Show Oth of N.-S. eq.'s }

x-component of vector form

$$\rho \left[\frac{\partial u}{\partial t} + \left(u \frac{\partial}{\partial x} + v \frac{\partial}{\partial y} + w \frac{\partial}{\partial z} \right) u \right] = -\frac{\partial p}{\partial x} + \mu \underbrace{\left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right)}_{\nabla^2 u}$$

$i=1$ for tensor form

$$\rho \left[\frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial x_j} \right] = -\frac{\partial p}{\partial x_i} + \mu \frac{\partial^2 u_i}{\partial x_j \partial x_j} \quad \nabla^2 u_i$$

$$\sum_{j=1}^3 u_j \frac{\partial u_i}{\partial x_j} = u_1 \frac{\partial u_i}{\partial x_1} + u_2 \frac{\partial u_i}{\partial x_2} + u_3 \frac{\partial u_i}{\partial x_3}$$

j is summation index, i is free (remaining) index

Einstein's summation convention: Σ is omitted
 (whenever an index appears twice, it is summed over)

Recitation 1 deals with Cartesian tensors,
 Kundu & Cohen, chapt. 2 also,
 and homework I.2.

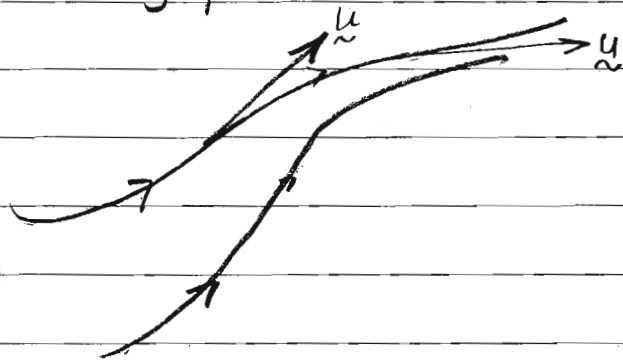
Kinematics:

Description of fluid particle motion, acceleration and deformation.

{ Shorter version exists }

Streamlines: tangents to the instantaneous velocity field

$$\frac{dx}{u} = \frac{dy}{v} = \frac{dz}{w}$$



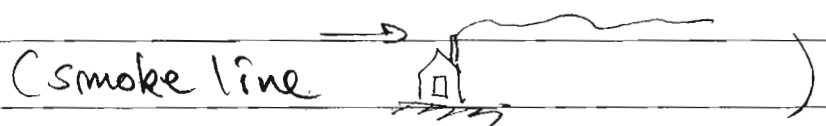
Pathlines: line following the path of a marked fluid particle

$$\frac{d\vec{x}_p(t)}{dt} = \vec{u}(\vec{x}_p(t), t)$$



{ For steady state flow these lines are the same }

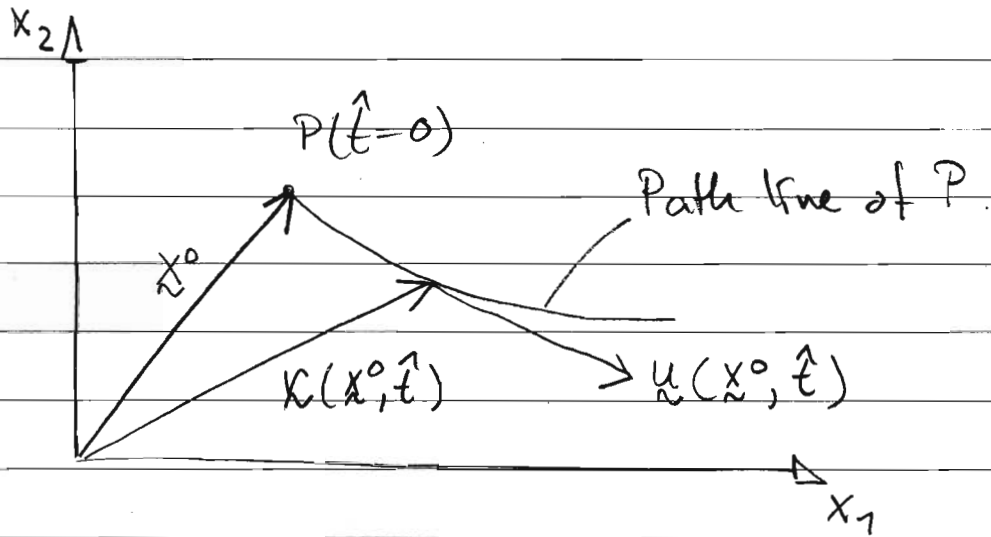
Streaklines: line connecting marked fluid particles that has passed a specific point in the flow



Timelines: line connecting neighbouring fluid particles at time t that were marked at t=0.

Lagrangian and Eulerian coordinates.

1. Lagrange coordinates (classical mechanics)



Every particle is marked and followed in the flow

Indep. variables : x_i^0 - initial position $i=1, 2, 3$

\hat{t} - time

$$r_i = r_i(x_1^0, x_2^0, x_3^0, \hat{t}) = r_i(x_i^0, \hat{t})$$

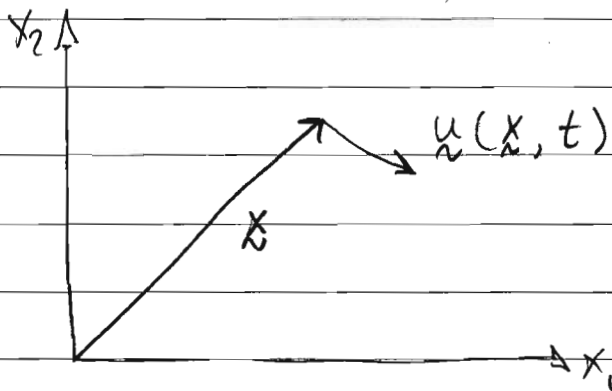
Velocity of particle $u_i = \frac{\partial r_i(x_i^0, \hat{t})}{\partial \hat{t}}$

Acceleration of part. $a_i = \frac{\partial^2 r_i}{\partial \hat{t}^2} = \frac{\partial u_i(x_i^0, \hat{t})}{\partial \hat{t}}$

When x_i^0 changes we consider new particles

Any flow variable $F(x_i^0, \hat{t})$

2. Euler coordinates



Consider fixed point in space. Fluid flows past point x at time t .

Independent variables : x_i - space coord. $i=1,2,3$

t - time

Fluid velocity at x_k and time t : $u_i(x_k, t)$

Flow variable $F(x_k, t)$

3. Relation between Lagrangian & Eulerian coordinates

x^0 defines a Lagrangian particle passing the point

$$\underline{x} = \underline{x}(x^0, \hat{t}) \quad \text{at the time } t = \hat{t}$$

The Eulerian velocity is $\underline{u}(x, t) = \frac{\partial}{\partial \hat{t}} \underline{x}(x^0, \hat{t})$

at $t = \hat{t}$, i.e. the velocity of the Lagrangian particle passing x at $t = \hat{t}$.

4. Material time derivative

$$F = F_L(\underline{x}^0, \hat{t}) = F_E(\underline{r}(\underline{x}^0, \hat{t}), t) \text{ at } t = \hat{t}$$

Rate of change of F for fluid particle

$$\frac{\partial}{\partial \hat{t}} F_L(\underline{x}^0, \hat{t})$$

Express this in Euler coordinates

$$\frac{\partial F_L}{\partial \hat{t}} = \frac{\partial F_E}{\partial x_i} \underbrace{\frac{\partial r_i}{\partial \hat{t}}}_{=u_i} + \frac{\partial F_E}{\partial t} \underbrace{\frac{\partial t}{\partial \hat{t}}}_{=1} = \frac{\partial F_E}{\partial t} + u_i \frac{\partial F_E}{\partial x_i}$$

$$\begin{aligned} \frac{\partial}{\partial \hat{t}} &= \frac{D}{Dt} \equiv \frac{\partial}{\partial t} + \underline{u} \cdot \nabla = \frac{\partial}{\partial t} + u_1 \frac{\partial}{\partial x_1} + u_2 \frac{\partial}{\partial x_2} + u_3 \frac{\partial}{\partial x_3} \\ &= \frac{\partial}{\partial t} + \underline{u} \cdot \nabla \end{aligned}$$

Material / substantial time derivative

$$\frac{D F}{D t} = \frac{\partial F}{\partial t} + \underline{u} \cdot \nabla F$$

Rate of change
for material
fluid particle

Rate of change
at fixed position
in space.

Advective rate of
change as particle
moves through spacial
gradients of F

Let \vec{u} be fluid velocity $\vec{u}(x, t)$

$$\frac{D}{Dt} \vec{u} = \frac{\partial \vec{u}}{\partial t} + \vec{u} \cdot \nabla \vec{u}$$

↑
Acceleration of material fluid particle.

Let \vec{x} be position vector \vec{x}

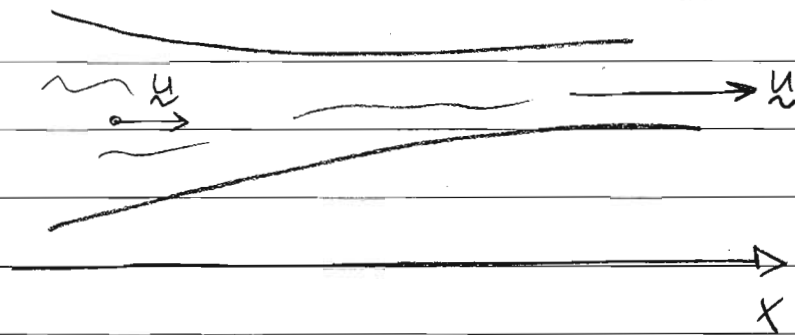
$$\frac{D}{Dt} \vec{x} = \frac{\partial \vec{x}}{\partial t} + \vec{u} \cdot \nabla \vec{x} = \underbrace{\frac{\partial \vec{x}}{\partial t}}_{=0} + u \frac{\partial \vec{x}}{\partial x} + v \frac{\partial \vec{x}}{\partial y} + w \frac{\partial \vec{x}}{\partial z}$$

$$= u \vec{e}_x + v \vec{e}_y + w \vec{e}_z = \vec{u}$$

$$\frac{D}{Dt} x_i = \frac{\partial x_i}{\partial t} + u_j \frac{\partial x_i}{\partial x_j} = \underbrace{\frac{\partial x_i}{\partial t}}_{=0} + u_j \delta_{ij} = u_i$$

Kronecker delta $\delta_{ij} = \begin{cases} 1 & ; \quad i=j \\ 0 & ; \quad i \neq j \end{cases}$; $\delta_{ij} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$

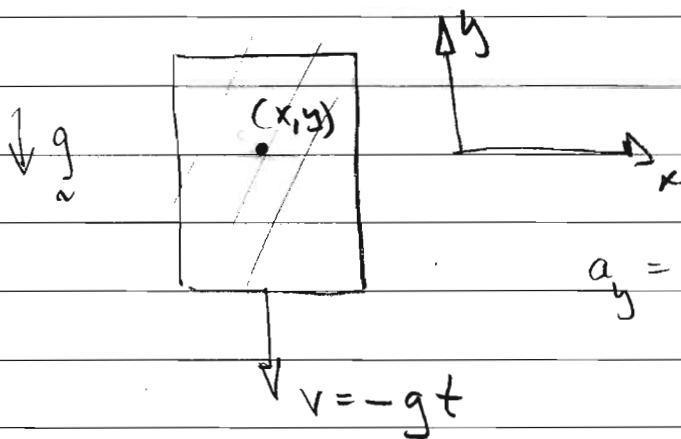
Acceleration of fluid particle in river



$$a_x = \underbrace{\frac{\partial u}{\partial t}}_{=0} + u \frac{\partial u}{\partial x} + \underbrace{v \frac{\partial u}{\partial y}}_{\text{assumed } 0} + \underbrace{w \frac{\partial u}{\partial z}}_{\text{assumed } 0} = \underbrace{u \frac{\partial u}{\partial x}}_{>0}$$

\therefore Acceleration of fluid element $\neq 0$ even if velocity at fixed x does not change with time.

Acceleration of falling can filled with fluid.



$$a_y = \frac{Dv}{Dt} = \frac{\partial v}{\partial t} + v \frac{\partial v}{\partial y} = -g - gt \cdot 0 = -g$$

\therefore No gradients of velocity field in bucket
 \Rightarrow advective term is zero.

See also homework I.1